



## Interweaving between collective and single-particle excitations in Sn-120 Andrea Idini

G. Potel, F. Barranco, E. Vigezzi, R.A. Broglia

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## **Two Degrees of Freedom**



#### Three Parameters

#### Hartree-Fock mean Field



 $m_k \approx 0.7m$ 

Pairing Interaction

 $v_{14}$ 

In <sup>120</sup>Sn with SLy4  $\Delta^{v_{14}} \approx 1.1 \text{ MeV}$  $\Delta^{exp} \approx 1.4 \text{ MeV}$  Particle Vibration Coupling Vertex

 $j_2$ 

 $\beta_{\lambda}$ 

2

0

E[MeV]

 $f(j_1, j_2, \lambda_{\nu}^{\pi})$ 

 $g(j_1, j_2, \lambda_{\nu}^{\pi}L)$ 

**RPA** 

<sup>120</sup>Sn

Exp.

# Nuclear Field Theory Approach



I<sup>st</sup> order

II<sup>nd</sup> order

Green's function can consistently Particle-Vibration Couplings to the <u>infinite order</u>:









#### **Pairing Properties**

# Two particle transfer cross section



## **Multiplet Splitting**

(keV

Elastic excitation of a quasiparticle state coupled to the core vibrations



 $(h_{11/2}\otimes 2^+)_{i^-}$ 400 11/2300 200 100 13/20 15/2100 7/2-200

#### **Electromagnetic Transition**



### Questioning the assumptions

Hartree-Fock mean Field



Pairing Interaction



Particle Vibration Coupling Vertex

















#### Questioning the assumptions



A neat description of Nuclear Structure

Interweaving single-particle (m<sub>k</sub> ≈ 0.7m) and collective (tuned to experiment) degrees of freedom, we can calculate several nuclear structure observables in open shell nuclei (pairing from realistic interactions) within 10% error.

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Interweaving of elementary modes of excitation in superfluid nuclei through particle-vibration coupling: Quantitative account of the variety of nuclear structure observables

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### **Multiplet Splitting**



$$\Sigma^{11}(a,\omega) = \sum_{b,\lambda^{\pi}_{\nu}} \int_{0}^{+\infty} \mathrm{d}\omega' \frac{V^{2}(a,b,\lambda^{\pi}_{\nu},\omega')}{\omega - \omega' - \hbar\omega_{\lambda^{\pi}_{\nu}}} + \frac{W^{2}(a,b,\lambda^{\pi}_{\nu},\omega')}{\omega + \omega' + \hbar\omega_{\lambda^{\pi}_{\nu}}},\tag{1.9}$$

$$\Sigma^{22}(a,\omega) = \sum_{b,\lambda^{\pi}_{\nu}} \int_{0}^{+\infty} \mathrm{d}\omega' \frac{V^{2}(a,b,\lambda^{\pi}_{\nu},\omega')}{\omega + \omega' + \hbar\omega_{\lambda^{\pi}_{\nu}}} + \frac{W^{2}(a,b,\lambda^{\pi}_{\nu},\omega')}{\omega - \omega' - \hbar\omega_{\lambda^{\pi}_{\nu}}},\tag{1.10}$$

$$\Sigma^{12}(a,\omega) = \sum_{b,\lambda^{\pi}_{\nu}} \int_{0}^{+\infty} \mathrm{d}\omega' V W(a,b,\lambda^{\pi}_{\nu},\omega') \left[ \frac{1}{\omega + \omega' + \hbar\omega_{\lambda^{\pi}_{\nu}}} - \frac{1}{\omega - \omega' - \hbar\omega_{\lambda^{\pi}_{\nu}}} \right]. \quad (1.11)$$



http://hdl.handle.net/2434/216315

$$\Sigma^{11}(a,\omega) = \sum_{b,\lambda^{\pi}_{\nu}} \int_{0}^{+\infty} \mathrm{d}\omega' \frac{V^{2}(a,b,\lambda^{\pi}_{\nu},\omega')}{\omega - \omega' - \hbar\omega_{\lambda^{\pi}_{\nu}}} + \frac{W^{2}(a,b,\lambda^{\pi}_{\nu},\omega')}{\omega + \omega' + \hbar\omega_{\lambda^{\pi}_{\nu}}},\tag{1.9}$$

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$$\Sigma^{12}(a,\omega) = \sum_{b,\lambda^{\pi}_{\nu}} \int_{0}^{+\infty} \mathrm{d}\omega' V W(a,b,\lambda^{\pi}_{\nu},\omega') \left[\frac{1}{\omega+\omega'+\hbar\omega_{\lambda^{\pi}_{\nu}}} - \frac{1}{\omega-\omega'-\hbar\omega_{\lambda^{\pi}_{\nu}}}\right]. \quad (1.11)$$

$$V^{2}(a, b, \lambda_{\nu}^{\pi}, \omega) = (f(a, b, \lambda_{\nu}^{\pi}) + g(a, b, \lambda_{\nu}^{\pi}))^{2} u^{2} a S^{+}(b, \omega) + (f(a, b, \lambda_{\nu}^{\pi}) - g(a, b, \lambda_{\nu}^{\pi}))^{2} v_{a}^{2} S^{-}(b, \omega) - 2(f(a, b, \lambda_{\nu}^{\pi}) + g(a, b, \lambda_{\nu}^{\pi}))(f(a, b, \lambda_{\nu}^{\pi}) - g(a, b, \lambda_{\nu}^{\pi}))u_{a} v_{a} \widetilde{S}(b, \omega)$$
(1.12)  
(1.13)

$$W^{2}(a, b, \lambda_{\nu}^{\pi}, \omega) = (f(a, b, \lambda_{\nu}^{\pi}) - g(a, b, \lambda_{\nu}^{\pi}))^{2} u_{a}^{2} S^{-}(b, \omega) + (f(a, b, \lambda_{\nu}^{\pi}) + g(a, b, \lambda_{\nu}^{\pi}))^{2} v_{a}^{2} S^{+}(b, \omega) - 2(f(a, b, \lambda_{\nu}^{\pi}) - g(a, b, \lambda_{\nu}^{\pi}))(f(a, b, \lambda_{\nu}^{\pi}) + g(a, b, \lambda_{\nu}^{\pi}))u_{a} v_{a} \widetilde{S}(b, \omega)$$
(1.14)

and

$$VW(a, b, \lambda_{\nu}^{\pi}, \omega) = (f(a, b, \lambda_{\nu}^{\pi}) + g(a, b, \lambda_{\nu}^{\pi}))^{2} u_{a} v_{a} S^{+}(b, \omega) - (f(a, b, \lambda_{\nu}^{\pi}) - g(a, b, \lambda_{\nu}^{\pi}))^{2} u_{a} v_{a} S^{-}(b, \omega) + (f(a, b, \lambda_{\nu}^{\pi}) + g(a, b, \lambda_{\nu}^{\pi}))(f(a, b, \lambda_{\nu}^{\pi}) - g(a, b, \lambda_{\nu}^{\pi}))(u_{a}^{2} - u_{a}^{2})\widetilde{S}(b, \omega)$$
(1.15)



$$\delta Vertex_{i}(a_{n}, b^{\mu}, \lambda_{\nu}^{\pi}, \omega) = \sum_{c,d,\lambda_{\nu'}^{\prime\pi'}} \frac{(-1)^{j_{c}+\lambda'-j_{a}}(-1)^{j_{d}+\lambda'-j_{b}}}{1+\delta_{\lambda_{\nu}^{\pi},\lambda_{\nu'}^{\prime\pi'}}} \langle (\lambda'j_{d})j_{b}\lambda; j_{a}|\lambda'(j_{d}\lambda)j_{c}; j_{a}\rangle$$

$$\sum_{\alpha,\beta} Vertex(a_{n}, c^{\alpha}, \lambda_{\nu'}^{\prime\pi'}) Vertex(c^{\alpha}, d^{\beta}, \lambda_{\nu}^{\pi}) Vertex(d^{\beta}, b^{\mu}, \lambda_{\nu'}^{\prime\pi'})$$

$$\frac{1}{\omega \mp (E_{c^{\alpha}} + \hbar\omega_{\lambda_{\nu'}^{\prime\pi'}})} \frac{1}{\omega \mp (E_{d^{\beta}} + \hbar\omega_{\lambda_{\nu'}^{\prime\pi'}} + \hbar\omega_{\lambda_{\nu}^{\pi}})} \tag{B.0.1}$$



**DWBA approximate** the entrance channel as a factorization of internal and relative coordinates, consider relative motion as Distorted Plain Wave, and calculate matrix element between this approximated *<intial| and |final>* state.

# 2-particle transfer DWBA



We need the structure information to calculate the correlation between the two transferred neucleons, so the probability of 1 neutron in the target and 1 in the ejectile, in the intermediate state

See G. Potel, A. Idini et al. Rep. Prog. Phys. 76 (106301)

#### one-particle transfer









